

# Interfacial Cracks in Adhesively Bonded Scarf Joints

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The interfacial crack behavior in adhesively bonded scarf joints is investigated. The study is conducted by using a method of analysis formulated on the basis of conservation laws of elasticity for nonhomogeneous solids and fundamental relationships in fracture mechanics of interfacial cracks. Solutions for interfacial crack-tip responses are determined for the joints with various adherend and adhesive properties, interface crack lengths, and adhesive-layer thicknesses. Also studied is the change of structural stiffness due to the presence of interfacial debonding. Of particular interests are the effects of geometric and material variables on the crack-tip response. Results are presented in terms of mixed-mode stress intensity factor solutions and energy release rates to illustrate the fundamental nature of fracture in adhesively bonded scarf joints.

## I. Introduction

APPLICATIONS of adhesive joining technology to advanced engineering structures such as aircraft construction have been of significant interest recently. Although adhesive joints have been employed in many secondary structural components, the reliable and effective use of adhesively bonded joints for primary load-bearing structures still remains in its infancy. The major factors responsible for this situation are the complex failure modes and mechanics presented in the joints. Among the various kinds of failure problems in adhesively bonded structures, interfacial cracking, sometimes also called debonding, often causes the greatest concern because it is one of the most commonly encountered failure mechanisms in the joints. Interfacial cracks are frequently observed to occur in fabrication and manufacturing processes such as trapped air bubbles or incomplete wetting between an adhesive and adherends. The cracks are also found to be initiated at geometric and material discontinuities such as at the edge of an adhesive-adherend interface. The presence of interfacial cracks in an adhesive joint can result in a progressive reduction of joint stiffness, the exposure of the interior of the composite to environmental attack, and the disintegration of the structure, which leads to the final fracture. Thus, understanding the basic nature of interfacial cracks and the ability of analyzing this critical problem are of paramount importance in the reliable design and future applications of the joints for primary load-bearing structures.

The scarf joint, which is considered to provide a higher strength than a lap-shear joint,<sup>1</sup> is studied because it is one of the most widely used configurations in transferring loads. Due to the complex structural geometry, the presence of a nonhomogeneous material system, the singular nature of debonding along the interface of dissimilar materials, and the unknown multiaxial stress state acting on the crack, the interfacial crack problem in adhesively bonded scarf joints is obviously very difficult. The very localized nature of the crack-tip response, as shown in Refs. 2-5, introduces the bond thickness as a characteristic dimension in the problem. Any simplifications that fail to include the thin but critically important adhesive layer would be unwise. The importance and complexities of interfacial debonding in adhesive joints have long been recognized<sup>6-9</sup>; however, research progress has

been relatively slow due to the aforementioned difficulties. In this paper a rigorous study of the mixed-mode interfacial crack problem in the composite structure is presented by using an approach recently developed by the authors.

The loading condition and the geometry of an adhesively bonded scarf joint under investigation are shown in Fig. 1a. The joint is composed of two high-stiffness and high-strength adherends bonded by a thin adhesive layer along a scarf angle  $\phi$  with respect to the applied nominal stress  $\sigma_\infty$  (Fig. 1b). The thin adhesive layer has elastic properties of  $E_1$  and  $\nu_1$ . The two adherends are made of high-modulus metals with material constants  $E_2$ ,  $\nu_2$  and  $E_3$ ,  $\nu_3$ . In this study, both the adhesive and the adherends are assumed to be linear elastic. The adhesive layer has a uniform thickness  $t$ ; the two adherends are considered to have the same uniform thickness  $d$ , and the joint, a total length of  $2L$ . A crack of length  $a$  is initiated at the traction-free edge and located along the interface between the adhesive and adherend. Except for the debonding region, the adherends and the adhesive are assumed to be perfectly bonded everywhere along bimaterial interfaces.

The objectives of this research are to: 1) develop an accurate and efficient method of analysis for examining the basic nature of the adhesive joint fracture problem; 2) investigate the complex failure modes and mechanics of interfacial cracks in adhesively bonded scarf joints; and 3) determine the effects of geometric and material variables on the crack behavior in the bonded composite structures.

The method of analysis is formulated on the basis of the recently developed conservation laws of elasticity for nonhomogeneous solids<sup>10-12</sup> and fundamental relationships in fracture mechanics of interfacial crack problems. The interfacial crack-tip response, characterized by inherently mixed-mode crack-tip stress intensity factors, is related to the conservation integrals derived in the paper. The conservation integrals can be evaluated conveniently by using properly introduced auxiliary solutions and field variables removed from the crack tip. The problem is then reduced to a pair of linear algebraic equations. Stress intensity and fracture energy release rate solutions are determined directly. One of the salient features of the present approach is that the stress intensity solution of an interfacial crack in a scarf joint can be determined accurately and efficiently by information extracted from the far field. The simplicity and accuracy of the current method make it particularly attractive for the present complex fracture problem. In the next section, formulation of conservation integrals for nonhomogeneous solids and their relationships with interfacial fracture mechanics parameters are presented. The solution procedure and the associated computational scheme are given in Sec. III. Solutions are shown in Sec. IV for the adhesively bonded scarf joints with

Presented as Paper 80-0775 at the AIAA/ASME/ASCE/AHS 21st Structures, Structural Dynamics and Materials Conference, Seattle, Wash., May 12-14, 1980; submitted May 28, 1980; revision received Feb. 18, 1981. Copyright © American Institute of Aeronautics and Astronautics, Inc., 1981. All rights reserved.

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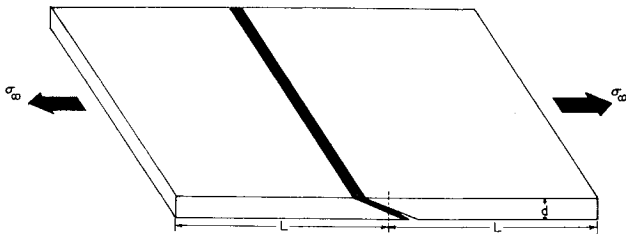


Fig. 1a Adhesively bonded scarf joint with an edge interfacial crack.

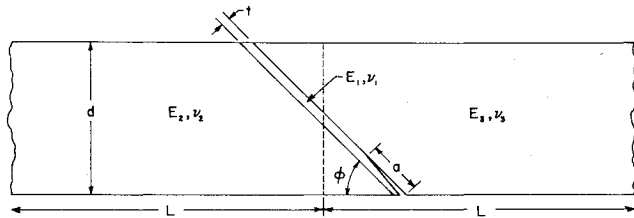


Fig. 1b Geometry and material constants of the joint and the crack.

various materials and geometric variables of the crack and adhesive layer. Many important features of the problem are revealed from the results, and a better understanding of the fundamental nature of the interfacial flaw behavior in adhesively bonded scarf joints is obtained. The method of analysis should be of practical use in a reliable design and analysis of adhesively bonded joints.

## II. Basic Formulation

For an infinitesimal deformation of a homogeneous elastic solid, a general form of the Eshelby-Rice conservation law<sup>13-16</sup> in a plane elasticity problem can be written as

$$J_i = \int_S (W n_i - \sigma_{jk} n_k u_{j,i}) ds = 0 \quad (1)$$

where  $W$  is the strain energy density,  $\sigma_{jk}$  the stress tensor,  $u_j$  the displacement vector, and  $n_i$  the outward unit normal of an arbitrary closed boundary  $S$  which encloses a portion of the continuum. The conservation integral  $J_i$  has recently been extended by Smelser and Gurtin<sup>10</sup> to a nonhomogeneous solid composed of dissimilar materials:

$$J_i = \int_S (W n_i - \sigma_{jk} n_k u_{j,i}) ds - \int_{\ell} ([W] \delta_{i2} - \sigma_{j2} [u_{j,i}]) ds = 0 \quad (2)$$

where  $\ell$  is a portion of the interface bounded by  $S$  as shown in Fig. 2a and the brackets denote the jump of a function across  $x_2 = 0$ , that is,

$$[W] = W(x_1, 0^+) - W(x_1, 0^-) \quad (3)$$

and

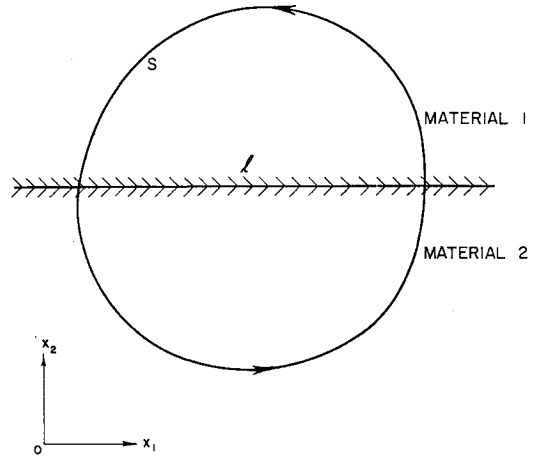
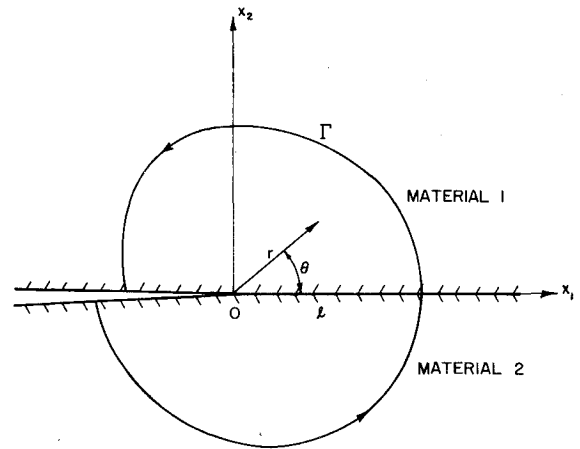
$$[u_{j,i}] = u_{j,i}(x_1, 0^+) - u_{j,i}(x_1, 0^-) \quad (4)$$

Continuity conditions of displacements and normal and shear stresses across the interface  $x_2 = 0$  require

$$[u_j] = 0 \quad \text{and} \quad [\sigma_{j2}] = 0 \quad (5)$$

Hence, the  $J_i$  component of the conservation integral in Eq. (2) can be simplified as

$$J_i = \int_S (W n_i - \sigma_{jk} n_k u_{j,i}) ds = 0 \quad (6)$$

Fig. 2a Closed integration path  $S$  for  $J_i$  in bimaterial composite.Fig. 2b Interfacial crack between dissimilar adherend and adhesive, and contour path  $\Gamma$  for  $J$  and  $M$  integrals.

which is the same as the analogous result for a single-phase material.

Consider a solid with a sharp crack located along the interface of two dissimilar materials, as shown in Fig. 2b. For the interfacial crack problem, the conservation integral of Eq. (6) can be shown to have the same form as the well-known  $J$  integral of fracture mechanics for monolithic solids without change, i.e.,

$$J \equiv J_i \{\Gamma\} = \int_{\Gamma} \left( W dx_2 - T_i \frac{\partial u_i}{\partial x_1} ds \right) \quad (7)$$

where  $\Gamma$  is an arbitrary path surrounding the crack tip, provided that the crack surfaces are free from traction and the interface is a straight line. It can also be shown that the  $J$  is equal to the energy release rate in the interface crack problem in a usual manner. Define the interfacial crack-tip stress intensity factors,  $K_I$  and  $K_{II}$ , by

$$K_I - iK_{II} = 2\sqrt{2\pi} e^{\beta\pi} \lim_{z \rightarrow 0} z^{1/2 + i\beta} \Phi_I(z) \quad (8)$$

where  $\Phi_I(z)$  is a complex potential in the well-known Kolosov-Muskhelishvili's formulation.<sup>8,17</sup> The  $\beta$  in Eq. (8) is the bimaterial constant

$$\beta = \frac{1}{2\pi} \ln \left[ \frac{(3-\nu_1)/[G_1(1+\nu_1)] + 1/G_2}{(3-\nu_2)/[G_2(1+\nu_2)] + 1/G_1} \right] \quad (9)$$

for a plane stress problem, where  $G_i$  and  $\nu_i$  are the shear modulus and Poisson's ratio of the  $i$ th material, respectively. For a plane strain case,  $\nu_i$  in Eq. (9) is replaced by  $\nu_i/(1-\nu_i)$ . One can easily show<sup>12,18</sup> that the  $J$  integral is related to stress intensity factors of an interfacial crack by

$$J = \mathcal{G} = \alpha(K_I^2 + K_{II}^2) \quad (10)$$

where  $\mathcal{G}$  is the total energy release rate defined by

$$\mathcal{G} = \mathcal{G}_I + \mathcal{G}_{II} \quad (11)$$

and  $\alpha$  has the form

$$\alpha = \begin{cases} \sum_{i=1}^2 \frac{1}{4G_i} (1 - \nu_i) & \text{(plane strain)} \\ \sum_{i=1}^2 \frac{1}{4G_i (1 + \nu_i)} & \text{(plane stress)} \end{cases} \quad (12)$$

Equation (10) was first given by Malyshev and Salganik.<sup>18</sup> The  $J$  integral alone does not provide adequate information for determining the individual values of  $K_I$  and  $K_{II}$  in an inherently mixed-mode interfacial fracture problem. However, if Eq. (10) is used in conjunction with conservation integrals for two independent equilibrium states of the nonhomogeneous elastic solid, this difficulty can be removed.

Now consider two independent equilibrium states of the elastically deformed bimaterial body. Associated field variables are denoted by superscripts 1 and 2. Superposition of the two equilibrium states leads to another equilibrium state, 0. The  $J$  integral for the superimposed state can be shown<sup>16</sup> to have the form

$$J^{(0)} = J^{(1)} + J^{(2)} + M^{(1,2)} \quad (13)$$

in which

$$M^{(1,2)} = \int_{\Gamma} \left( W^{(1,2)} dx_2 - \left[ T_I^{(1)} \frac{\partial u_I^{(2)}}{\partial x_I} + T_I^{(2)} \frac{\partial u_I^{(1)}}{\partial x_I} \right] ds \right) \quad (14)$$

where  $W^{(1,2)}$  is the mutual potential energy density of the nonhomogeneous body defined by

$$W^{(1,2)} = C_{ijkl} u_{i,j}^{(1)} u_{k,l}^{(2)} = C_{ijkl} u_{i,j}^{(2)} u_{k,l}^{(1)} \quad (15)$$

Recalling the  $J$ - $K$  relationship in Eq. (10) and the orthogonality of stress intensity factors, one can express the  $J$  integral for the new state by

$$J^{(0)} = \alpha([K_I^{(1)} + K_I^{(2)}]^2 + [K_{II}^{(1)} + K_{II}^{(2)}]^2) \quad (16)$$

More explicitly, Eq. (16) can be written as

$$J^{(0)} = J^{(1)} + J^{(2)} + 2\alpha(K_I^{(1)} K_I^{(2)} + K_{II}^{(1)} K_{II}^{(2)}) \quad (17)$$

A direct comparison between Eqs. (13) and (17) leads to

$$M^{(1,2)} = 2\alpha[K_I^{(1)} K_I^{(2)} + K_{II}^{(1)} K_{II}^{(2)}] \quad (18)$$

The  $M$  integral in Eqs. (14) and (18) deals with interaction terms only, and will be used directly in solving the interfacial crack problem. The integral is related to the details of the stresses and deformation at the crack tip [i.e.,  $K_I$  and  $K_{II}$  in Eq. (18)], but yet can be evaluated in a region away from the

crack tip [i.e., the integral in Eq. (14)], where such a calculation can be carried out with greater accuracy and convenience than near the crack tip.

### III. Solution Procedure

Equation (14) together with Eq. (18) provides, in fact, sufficient information for determining the stress intensity solutions  $K_I^{(1)}$  and  $K_{II}^{(1)}$  for the mixed-mode interfacial crack problem, when proper auxiliary solutions are introduced. Denote the first auxiliary solution by a superscript 2a. For an interfacial crack in a body of dissimilar materials subjected to mode I deformation only, i.e.,

$$K_I^{(2a)} = I \quad \text{and} \quad K_{II}^{(2a)} = 0 \quad (19)$$

Equation (18) can be simplified as

$$M^{(1,2a)} = 2\alpha K_I^{(1)} \quad (20)$$

where  $M^{(1,2a)}$  has the form

$$M^{(1,2a)} = \int_{\Gamma} \left( C_{ijkl} u_{i,j}^{(1)} u_{k,l}^{(2a)} dx_2 - \left[ T_I^{(1)} \frac{\partial u_I^{(2a)}}{\partial x_I} + T_I^{(2a)} \frac{\partial u_I^{(1)}}{\partial x_I} \right] ds \right) \quad (21)$$

and  $T_I^{(1)}$  and  $u_I^{(1)}$  in Eq. (21) can be determined along an arbitrarily selected integration path,  $\Gamma$ . Any convenient method such as the commonly used finite-element analysis, which can provide an accurate far-field solution, is suitable for this purpose. For a plane strain crack problem with the loading of Eq. (19),  $T_I^{(2a)}$  and  $u_I^{(2a)}$  are derivable by the Kolosov-Muskhelishvili formulation and have the forms:

$$u_I^{(2a)} = \sqrt{\frac{r}{2\pi}} f_I^{(1)}(\ln r, \theta; G_1, \nu_1, G_2, \nu_2) \quad (22)$$

and

$$T_I^{(2a)} = \sigma_{ij}^{(2a)} n_j \quad (23)$$

with

$$\sigma_{ij}^{(2a)} = \frac{1}{\sqrt{2\pi r}} g_{ij}^{(1)}(\ln r, \theta; G_1, \nu_1, G_2, \nu_2) + O(1) \quad (24)$$

Exact forms of the functions  $f_I^{(1)}$  and  $g_{ij}^{(1)}$  can be found in Refs. 8, 18, and 19.

The second auxiliary solution is denoted by the superscript 2b for a pure mode II deformation of the nonhomogeneous solid with

$$K_I^{(2b)} = 0 \quad \text{and} \quad K_{II}^{(2b)} = I \quad (25)$$

Substituting Eq. (25) into Eq. (18), one obtains the following relationship

$$M^{(1,2b)} = 2\alpha K_{II}^{(1)} \quad (26)$$

where

$$M^{(1,2b)} = \int_{\Gamma} \left( C_{ijkl} u_{i,j}^{(1)} u_{k,l}^{(2b)} dx_2 - \left[ T_I^{(1)} \frac{\partial u_I^{(2b)}}{\partial x_I} + T_I^{(2b)} \frac{\partial u_I^{(1)}}{\partial x_I} \right] ds \right) \quad (27)$$

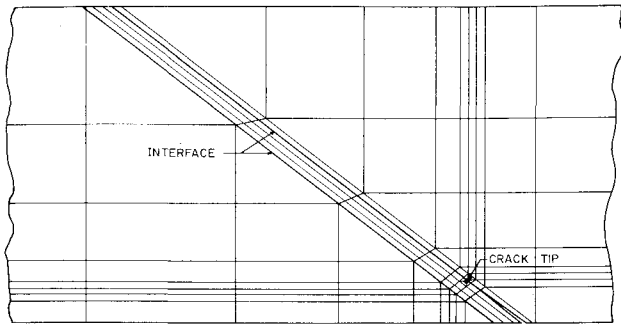


Fig. 3a Schematic finite-element mesh configuration around interfacial crack and adhesive layer.

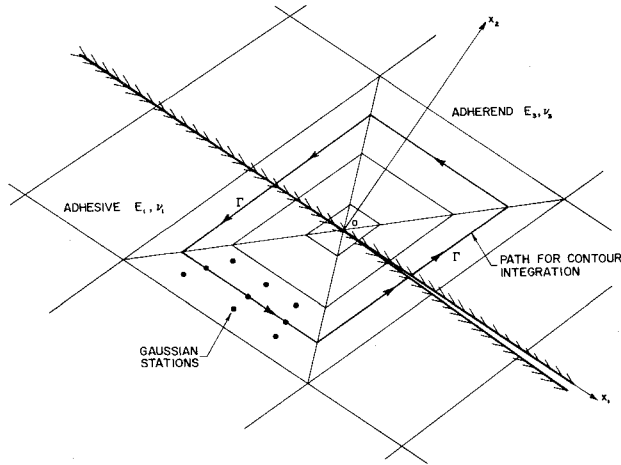


Fig. 3b Path for contour integration and finite-element mesh arrangement in crack-tip region.

The auxiliary solutions  $u_i^{(2b)}$  and  $T_i^{(2b)}$  have the form

$$u_i^{(2b)} = \sqrt{\frac{r}{2\pi}} f_i^{(II)}(\ln r, \theta; G_1, \nu_1, G_2, \nu_2) \quad (28)$$

and

$$T_i^{(2b)} = \sigma_{ij}^{(2b)} n_j \quad (29)$$

where

$$\sigma_{ij}^{(2b)} = \frac{I}{\sqrt{2\pi r}} g_{ij}^{(II)}(\ln r, \theta; G_1, \nu_1, G_2, \nu_2) + O(1) \quad (30)$$

It is important to note that the auxiliary solutions,  $u_i^{(2a)}$ ,  $T_i^{(2a)}$  and  $u_i^{(2b)}$ ,  $T_i^{(2b)}$ , are independent of the particular boundary-value problem being posed. Therefore, they can be determined independently by any convenient analytical method once and for all. In solving for  $K_I^{(I)}$  and  $K_{II}^{(I)}$  from Eqs. (20), (21), (26), and (27), one must evaluate the integrals  $M^{(I,2a)}$  and  $M^{(I,2b)}$  accurately and explicitly. For a given joint geometry, loading condition, and material constants, the evaluation can be made by integrating Eqs. (21) and (27) along a properly selected contour in the far field in order to avoid crack-tip complications. A numerical method, using a conventional finite-element approach, is currently employed to calculate  $T_i^{(I)}$  and  $u_i^{(I)}$ . The  $M$  integral is then evaluated by a numerical integration scheme using the second-order Gaussian quadrature along a properly selected contour path  $\Gamma$  away from the crack tip.

#### IV. Numerical Results and Discussions

The solution procedure outlined in Sec. III has been programmed for studying general two-dimensional interfacial crack problems in adhesively bonded joints. Evaluation of the

conservation contour integrals is conducted in conjunction with a finite-element method by using eight-node isoparametric elements. For illustrative purposes, a typical finite-element mesh for a scarf joint consisting of 182 elements used in actual computation is shown schematically in Fig. 3a in the neighborhood of the adhesive bond region containing the interfacial crack. The detailed element arrangement around the crack tip and the path selected for contour integration are given in Fig. 3b. The path independency of the conservation integrals,  $J$  and  $M$ , has been checked numerically by using different paths of integration. Due to the inherently singular nature of the problem and the cumulative nature of errors introduced by the numerical analysis, proper selection of the path for integration removed enough from the crack tip is essential to insure the path independency of the contour integrals. Numerical details regarding the choice of the path are discussed in several related papers,<sup>11,20,21</sup> and are not repeated here. Solutions are obtained for adhesive scarf joints with various materials and geometric variables. Accuracy and convergence of the results are affected by several unusual features of the problem and the method of analysis due to the singular nature of the crack and inherent approximation in the numerical evaluation of the integrals. Assessments of the solution accuracy are made by examining relevant problems for which unquestionably correct and exact solutions are available in the literature. Excellent agreement is observed between current results and reference solutions. Details of the study are reported elsewhere.<sup>11,20</sup> The current results, based on the optimum mesh arrangement and the integration contour selected, have an accuracy within approximately 3% deviation from the reference solutions.

Now consider an adhesively bonded scarf joint containing an interfacial crack subjected to a unit in-plane nominal stress,  $\sigma_\infty$  [i.e.,  $\sigma_\infty = 6.895 \times 10^3$  Pa (1 psi)], as shown in Figs. 1a and 1b. The adherends in the composite structure have a uniform thickness  $d = 1.27$  cm (0.5 in.) and a total length of  $2L = 25.4$  cm (10 in.). For simplicity without loss of generality, the scarf angle  $\phi$  of the joint is taken as 45 deg for illustrative purposes ( $\phi$  is generally smaller than 45 deg in most practical cases). The adherends are structural aluminum with elastic constants typical for aircraft construction

$$E_2 = E_3 = 6.895 \times 10^4 \text{ MPa } (10 \times 10^6 \text{ psi})$$

and

$$\nu_2 = \nu_3 = 0.33$$

Effects of geometric and material variables of the thin adhesive layer and the interfacial crack on the fracture behavior of the joints are the subject under current investigation. The primary objectives of this section are to examine the complex failure modes and mechanics associated with an interfacial crack and to assess the influence of problem variables on the crack-tip behavior in adhesively bonded scarf joints. Mixed-mode stress intensity factor solutions and associated energy release rates, which characterize the stresses and deformation in the crack-tip region, are evaluated for various cases. The reduction of structural stiffness of the joint due to the presence of an edge interfacial crack is also studied to illustrate the significance of the defect on structural integrity. Of particular interest is the crack-tip response in the joints with different bond thicknesses; quantitative information on this important subject is determined. In the following numerical examples, the plane strain condition is assumed for all of the interfacial crack problems.

##### A. Effect of Adhesive Material Property

Effects of different kinds of adhesive on the crack-tip response in the joints are examined by considering adhesives

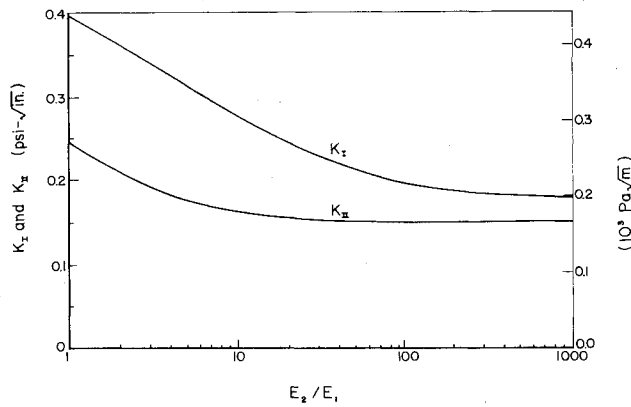


Fig. 4 Effect of adhesive modulus on mixed-mode interfacial crack-tip stress intensity factors ( $E_2 = E_3 = 6.895 \times 10^4$  MPa ( $10 \times 10^6$  psi)).

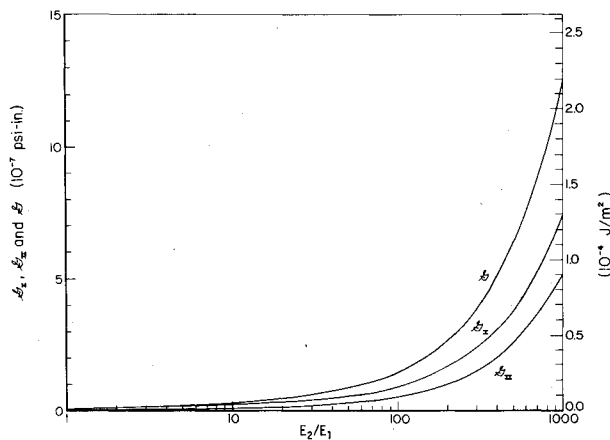


Fig. 5 Effect of adhesive modulus on energy release rates in debonded scarf joint ( $E_2 = E_3 = 6.895 \times 10^4$  MPa ( $10 \times 10^6$  psi)).

with various material constants,  $E_1$ 's or  $E_2/E_1$ 's (with given adherends and  $\nu_1 = 0.3$ ). For the purpose of generality, adhesives with a broad range of Young's modulus covering three decades on a logarithmic scale are studied. The cases with  $E_2/E_1 = 20 \sim 40$  are typical for aluminum-epoxy composite structures. Higher  $E_2/E_1$  values correspond to the joints with less rigid adhesives or subjected to a "hot and wet" environment. Crack-tip stress intensity factor solutions and associated energy release rates are shown in Figs. 4 and 5 as functions of the modulus ratio for the joints with  $t/d = 0.05$  and  $a/d = 0.2$ . Failure modes and mechanics associated with the interfacially cracked joints are revealed. Obviously, the interfacial crack experiences a mixed-mode fracture even under uniaxial nominal loading, because  $K_I$  and  $K_{II}$  are of the same order of magnitude for all of the cases studied.

Both  $K_I$  and  $K_{II}$  are observed to decrease as the modulus of the adhesive decreases. The difference between  $K_I$  and  $K_{II}$  decreases with  $E_1$  also. The mode I stress intensity factor is always higher than that of the mode II in the entire range of  $E_2/E_1$  studied; thus, the dominant mode of failure is an opening one in the joints with the scarf angle considered. The shearing mode stress intensity factor remains relatively unchanged for most practical adhesive scarf joints with  $E_2/E_1$  greater than 20. In the case of  $E_2/E_1 = 1$ , the crack is cohesive, and the nature of the crack-tip singularity changes. In this case, stress intensity solutions,  $K_I$  and  $K_{II}$ , are found to be consistent with the results given in Ref. 22. The energy release rate of each individual fracture mode,  $G_I$  and  $G_{II}$ , and the total energy release rate,  $G$ , are evaluated. For the joints with a given adherend material, alternation of adhesive properties results in a change of  $\alpha$  in Eqs. (10-12). The reduction of  $E_1$  increases the  $\alpha$  value which significantly influences the  $G_I$ ,  $G_{II}$  and  $G$ . Both  $G_I$  and  $G_{II}$  increase rapidly

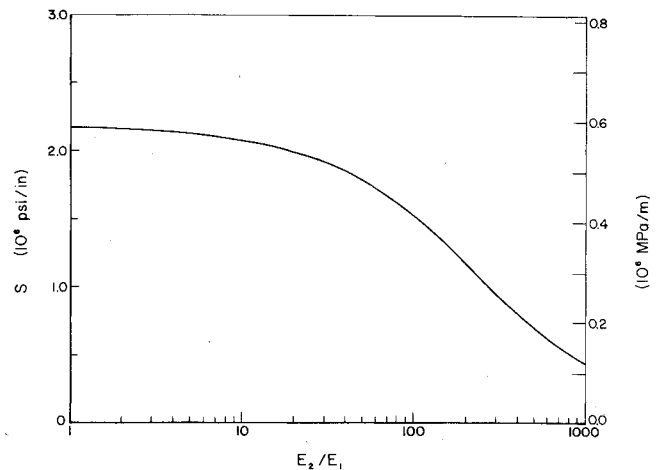


Fig. 6 Variation of structural stiffness,  $S$ , of debonded adhesive scarf joints with different  $E_2/E_1$ 's ( $E_2 = E_3 = 6.895 \times 10^4$  MPa ( $10 \times 10^6$  psi)).

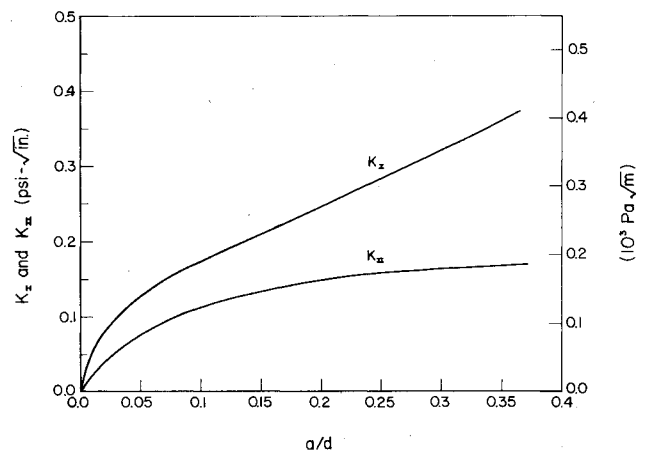


Fig. 7 Effect of interface crack length,  $a/d$  on, mixed-mode crack-tip stress intensity factors  $K_I$  and  $K_{II}$ .

with  $E_2/E_1$ , especially as  $E_2/E_1 > 20$ . The rapid increase of  $G_I$ ,  $G_{II}$  and  $G$  corresponding to the decrease of  $K_I$  and  $K_{II}$  is noted to be a phenomenon uniquely associated with the interfacial crack problem, and is not observed in a homogeneous crack problem in general. In the cases of very large values of  $E_2/E_1$ , the adhesive may become incompressible with a Poisson's ratio approaching 0.5. The incompressibility of the soft adhesive is not considered in the present study.

The influence of adhesive material properties on structural integrity of debonded scarf joints is shown in Fig. 6 in which a reduction of structural stiffness,  $S$ , is observed (computed on the basis of the applied stress and resulting displacements at the ends of the joint). The joint stiffness decreases monotonically with the adhesive modulus. The stiffness reduction is found especially significant for the joints with  $E_2/E_1$  greater than 20. For instance, an approximately 400% decrease in the joint structural stiffness is observed, as  $E_2/E_1$  increases from 20 to 1000.

#### B. Influence of Interfacial Crack (Debonding) Length

The interfacial crack length is of particular interest because of its influences on failure mechanics, joint integrity, and overall performance of the bonded structure. The joints studied here have the same geometry as those examined in the previous section [i.e.,  $2L = 25.4$  cm (10 in.),  $d = 1.27$  cm (0.5 in.),  $\phi = 45$  deg,  $t/d = 0.05$ ]. The adherends are the same aluminum, and the adhesive is a structural epoxy with material constants  $E_1 = 3.448 \times 10^3$  MPa ( $0.5 \times 10^6$  psi) and

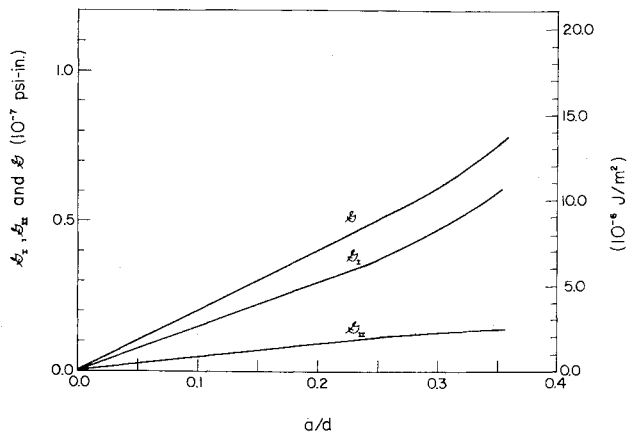


Fig. 8 Effect of interface crack length,  $a/d$ , on energy release rates,  $G_I$ ,  $G_{II}$  and  $G$  in debonded scarf joint.

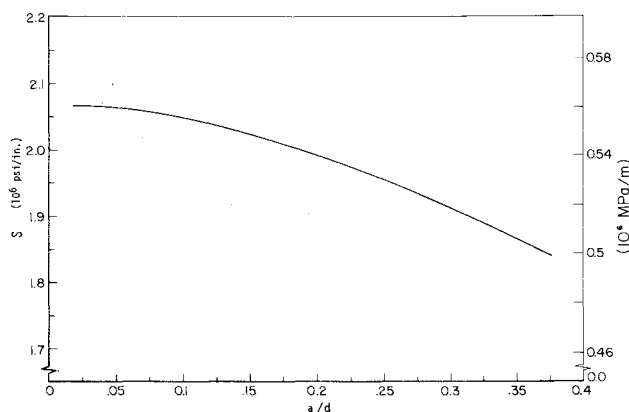


Fig. 9 Reduction of structural stiffness,  $S$ , due to the presence of interface crack,  $a/d$ , in debonded scarf joint.

$\nu_I = 0.3$ . The scarf joints with interfacial crack lengths,  $a/d$ , ranging from 0.05 to 0.4 are studied in detail. Effects of  $a/d$  on mixed-mode crack-tip stress intensity factors are shown in Fig. 7. The opening mode stress intensity factor  $K_I$  increases monotonically with the debonding length. The shearing mode stress intensity  $K_{II}$  changes with the interfacial crack length in a manner similar to  $K_I$  but at a lower rate. As the debonding length exceeds  $0.2d$ ,  $K_{II}$  becomes less sensitive to the change of  $a/d$ . The value of  $K_I$  is found increasingly larger than that of  $K_{II}$  as  $a/d$  increases. Hence, the larger the debonding length, the more dominant the opening mode fracture in the joint. Corresponding energy release rates in the joints with various interfacial crack lengths are reported in Fig. 8. As anticipated,  $G_I$ ,  $G_{II}$  and  $G$  change with the crack length in an approximately linear manner. Slight deviation from linearity occurs as  $a/d$  becomes large. The mode II energy release rate  $G_{II}$  is observed to increase with a smaller gradient as compared with that of  $G_I$ . The associated structural stiffness change of the joint is evaluated also and shown in Fig. 9. As one expects, increasing the interfacial crack length leads to a gradual reduction of joint stiffness. For example, a change of debonding length  $a/d$  from 0.1 to 0.4 causes a 15% decrease in structural stiffness.

### C. Effect of Adhesive-Layer Thickness

Effects of the adhesive-layer thickness on fracture of bonded joints have long been recognized. Bascom et al.,<sup>5</sup> for example, showed experimentally that fracture energy release rates were related closely to the bond thickness, indicating the criticality of the adhesive thickness in controlling the failure of structural joints. Erdogan et al.<sup>2,23</sup> conducted analytical studies on crack-tip stresses in debonded composite laminates and reported unique features of this geometric variable in the

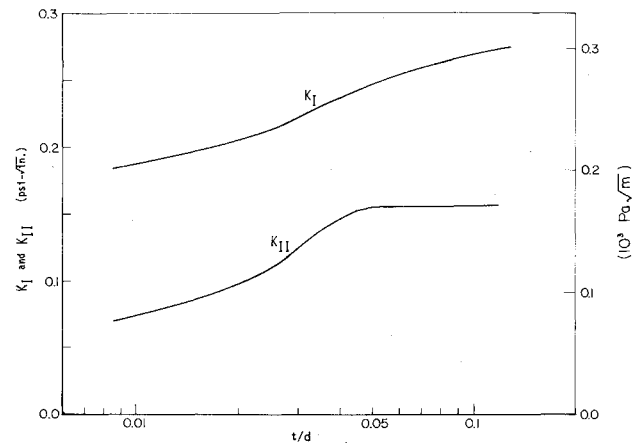


Fig. 10 Effect of adhesive layer thickness,  $t/d$ , on interface crack-tip stress intensity factors  $K_I$  and  $K_{II}$ .

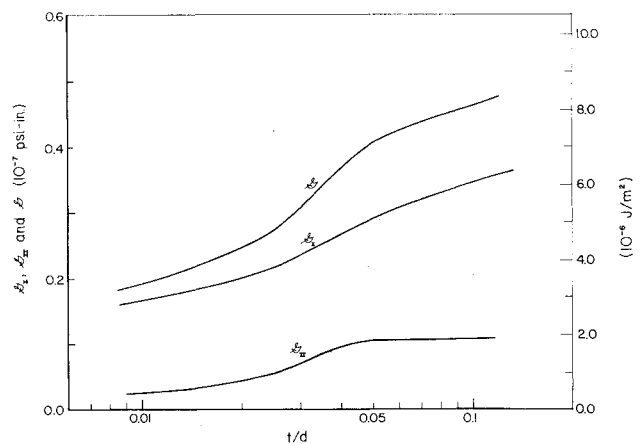


Fig. 11 Effect of adhesive layer thickness,  $t/d$ , on energy release rates,  $G_I$ ,  $G_{II}$  and  $G$  in debonded scarf joint.

joint fracture behavior. Wang et al.<sup>24,25</sup> examined center-of-the-bond cohesive cracks and showed fundamental characteristics of thickness effects of the adhesive layer on fracture of the joints. In this portion of the study, scarf joints containing an interfacial crack of  $a/d = 0.2$  with various adhesive-layer thicknesses,  $t/d$ , ranging from 0.01 to 0.1 [i.e.,  $t = 0.127 \sim 1.27$  mm (0.005–0.05 in.)] are considered. Other geometric variables and materials properties are identical to those used in previous sections.

Relationships among  $K_I$ ,  $K_{II}$ , and the bond thickness are shown in Fig. 10. The opening mode stress intensity solution is found to increase initially with the adhesive-layer thickness. As  $t/d$  becomes very large,  $K_I$  and  $K_{II}$  must approach an asymptotic value for the interfacial crack between two dissimilar materials, i.e., the adhesive  $E_I$  and the adherend  $E_3$ . The shearing mode stress intensity factor increases with  $t/d$  first, and then becomes less sensitive to the change of the bond thickness as  $t$  exceeds a certain value (i.e.,  $t/d = 0.05$  in this case). The opening mode stress intensity factor is always larger than that of the shearing mode in all cases studied; thus, the mode I failure becomes increasingly critical as the adhesive thickness increases. Associated energy release rates for the joints with various  $t/d$ 's are given in Fig. 11. With a given debonding length, the decrease of the adhesive-layer thickness leads to an increase of fracture resistance in the bonded joints (i.e., a reduction in the total energy release rate and in the maximum cleavage stress). The change of fracture resistance with the bond thickness is of critical importance in the fracture design of adhesively bonded structures. This finding is consistent with the results reported by Erdogan.<sup>23</sup> However, variations of interfacial crack-tip stress intensity

factors and associated energy release rates with  $t/d$  in the current problem are in contrast to the solutions obtained for center-of-the-bond cohesive crack problems in which  $K_I$ 's and  $G_I$ 's are found to be almost independent of the adhesive-layer thickness. This situation is expected as the associated stress singularities are quite different between the two cases,<sup>24</sup> and the stress intensity factors of an interfacial crack are defined differently from those of a cohesive one.

### V. Summary and Conclusions

Interfacial cracks in adhesively bonded scarf joints are studied with a method of analysis based on conservation laws in elasticity for nonhomogeneous solids and fundamental relationships of fracture mechanics for interfacial crack problems. The current approach is shown to provide an effective and efficient means to examine the complex interfacial crack behavior and the performance of debonded joints. Solutions are obtained for scarf joints containing an edge interfacial crack with various adhesive types, debonding lengths, and adhesive-layer thicknesses. Fracture parameters such as mixed-mode stress intensity factors and associated energy release rates describing the crack-tip stresses and deformation are determined. Influences of material and geometric variables on the integrity of the bonded structures are also studied.

Based on the results obtained in the study, the following conclusions can be drawn:

1) The current method of analysis is superior to other approaches in its operational simplicity and solution accuracy for studying mixed-mode interfacial crack problems in adhesive joints with complex geometries.

2) The change of adhesive material properties significantly influences the interfacial crack-tip response and joint performance. In the scarf joints studied, the opening mode stress intensity factor is always greater than that of the shearing mode, and the difference between  $K_I$  and  $K_{II}$  decreases with the increase of  $E_2/E_1$ . The associated energy release rates increase rapidly with the decrease of crack-tip stress intensity factors as  $E_2/E_1$  increases—a phenomenon unique to the interfacial crack behavior in a nonhomogeneous solid and not observed in homogeneous crack problems in general. The structural stiffness of a debonded joint decreases rapidly as  $E_2/E_1$  becomes greater than 20.

3) Increasing the interfacial crack length in scarf joints raises  $K_I$  and  $K_{II}$  simultaneously. The magnitude and the rate of change of  $K_I$  are considerably larger than those of  $K_{II}$ . As the crack length increases,  $K_{II}$  becomes gradually insensitive to  $a/d$ , but  $K_I$  increases monotonically. Energy release rates associated with the interfacial crack increase approximately linearly with the crack length, but deviation from linearity starts as the crack length  $a/d$  exceeds 0.3. The structural stiffness of the joint decreases appreciably with  $a/d$ .

4) Effects of the adhesive-layer thickness on the interfacial crack response are significant. Both  $K_I$  and  $G_I$  increase monotonically with the change of the adhesive thickness initially, and approach their asymptotic values, which are the mode I stress intensity and energy release rate for an interfacial crack in two dissimilar materials, as  $t/d$  becomes very large. But  $K_{II}$  and  $G_{II}$  become relatively insensitive to the increase of  $t/d$  as the adhesive layer exceeds a certain thickness. The increase of the adhesive bond thickness leads to a decrease of fracture resistance of the joint in general.

### Acknowledgments

The work reported in this paper was supported in part by the Office of Naval Research through Contract N00014-79-C-0579. Computation was carried out in the Digital Computer Laboratory of the University of Illinois at Urbana-Champaign. The authors are grateful to Dr. N. Perrone of ONR for his encouragement and cooperation during the course of this study. Fruitful discussion with Dr. D. L. Hunston of the Naval Research Laboratory, Washington, D. C., is also gratefully acknowledged.

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